

# Regularity and mass conservation for discrete coagulation-fragmentation equations with diffusion

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## Abstract

We present a new a-priori estimate for discrete coagulation-fragmentation systems with size-dependent diffusion within a bounded, regular domain confined by homogeneous Neumann boundary conditions. Following from a duality argument, this a-priori estimate provides a global  $L^2$  bound on the mass density and was previously used, for instance, in the context of reaction-diffusion equations.

In this paper we demonstrate two lines of applications for such an estimate: On the one hand, it enables to simplify parts of the known existence theory and allows to show existence of solutions for generalised models involving collision-induced, quadratic fragmentation terms for which the previous existence theory seems difficult to apply. On the other hand and most prominently, it proves mass conservation (and thus the absence of gelation) for almost all the coagulation coefficients for which mass conservation is known to hold true in the space homogeneous case.

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# 1 Introduction

We consider the time evolution of a physical system where a set of particles can aggregate into groups of two or more, called *clusters*, and where these clusters can diffuse in space with a diffusion constant which depends on their size. If we represent space by an open bounded set  $\Omega \subseteq \mathbb{R}^N$  with regular boundary, the initial-boundary problem for the concentrations  $c_i = c_i(t, x) \geq 0$  of clusters with integer size  $i \geq 1$  at position  $x \in \Omega$  and time  $t \geq 0$  is given by the discrete coagulation-fragmentation system of equations with spatial diffusion and homogeneous Neumann boundary conditions :

$$\partial_t c_i - d_i \Delta_x c_i = Q_i + F_i \quad \text{for } x \in \Omega, t \geq 0, i \in \mathbb{N}^*, \quad (1a)$$

$$\nabla_x c_i \cdot n = 0 \quad \text{for } x \in \partial\Omega, t \geq 0, i \in \mathbb{N}^*, \quad (1b)$$

$$c_i(0, x) = c_i^0(x) \quad \text{for } x \in \Omega, i \in \mathbb{N}^*, \quad (1c)$$

where  $n = n(x)$  represents a unit normal vector at a point  $x \in \partial\Omega$ ,  $d_i$  is the diffusion constant for clusters of size  $i$ , and

$$Q_i \equiv Q_i[c] := Q_i^+ - Q_i^- := \frac{1}{2} \sum_{j=1}^{i-1} a_{i-j,j} c_{i-j} c_j - \sum_{j=1}^{\infty} a_{i,j} c_i c_j, \quad (2)$$

$$F_i \equiv F_i[c] := F_i^+ - F_i^- := \sum_{j=1}^{\infty} B_{i+j} \beta_{i+j,i} c_{i+j} - B_i c_i.$$

The parameters  $B_i$ ,  $\beta_{i,j}$  and  $a_{i,j}$ , for integers  $i, j \geq 0$ , represent the total rate  $B_i$  of fragmentation of clusters of size  $i$ , the average number  $\beta_{i,j}$  of clusters of size  $j$  produced due to fragmentation of a cluster of size  $i$ , and the coagulation rate  $a_{i,j}$  of clusters of size  $i$  with clusters of size  $j$ . We refer to these parameters as *the coefficients* of the system of equations. They represent rates, so they are always nonnegative; single particles do not fragment further, and mass should be conserved when a cluster fragments into smaller pieces, so one always imposes

$$a_{i,j} = a_{j,i} \geq 0, \quad \beta_{i,j} \geq 0, \quad (i, j \in \mathbb{N}^*), \quad (3a)$$

$$B_1 = 0, \quad B_i \geq 0, \quad (i \in \mathbb{N}^*), \quad (3b)$$

$$i = \sum_{j=1}^{i-1} j \beta_{i,j}, \quad (i \in \mathbb{N}, i \geq 2). \quad (3c)$$

In fact, the last condition (3c) implies the conservation of the total mass  $\int_{\Omega} \sum_{i=1}^{\infty} i c_i dx$ , which becomes obvious from the following formal *fundamental identity* or *weak formulation* of the coagulation and fragmentation operators:

Consider a sequence of nonnegative numbers  $\{c_i\}$ , and define  $Q_i, F_i$  as in eqs. (2), then, for any sequence of numbers  $\varphi_i$ ,

$$\begin{aligned} \sum_{i=1}^{\infty} \varphi_i Q_i &= \frac{1}{2} \sum_{i=1}^{\infty} \sum_{j=1}^{\infty} a_{i,j} c_i c_j (\varphi_{i+j} - \varphi_i - \varphi_j), \\ \sum_{i=1}^{\infty} \varphi_i F_i &= - \sum_{i=2}^{\infty} B_i c_i \left( \varphi_i - \sum_{j=1}^{i-1} \beta_{i,j} \varphi_j \right). \end{aligned} \quad (4)$$

As a (still formal) consequence for solutions  $\{c_i\}$  of (1) – (2), one can calculate the time derivative of the integral of the moment  $\sum \varphi_i c_i$  to obtain

$$\frac{d}{dt} \int_{\Omega} \sum_{i=1}^{\infty} \varphi_i c_i = \int_{\Omega} \sum_{i=1}^{\infty} \varphi_i (Q_i + F_i), \quad (5)$$

since the integral of the diffusion part vanishes due to the homogeneous Neumann boundary condition. By choosing  $\varphi_i := i$  above and thanks to (3c), we have  $\sum_{i=1}^{\infty} i Q_i = \sum_{i=1}^{\infty} i F_i = 0$ , and the total mass is formally conserved :

$$\|\rho(t, \cdot)\|_{L^1} = \int_{\Omega} \sum_{i=1}^{\infty} i c_i(t, x) dx = \int_{\Omega} \sum_{i=1}^{\infty} i c_i^0(x) dx = \|\rho^0\|_{L^1} \quad (t \geq 0). \quad (6)$$

Our main aim in this work is to provide some new bounds on the regularity of weak solutions for system (1) – (2) by means of techniques developed in the context of reaction-diffusion equations [9, 16, 17], and to give three applications to those bounds, the main one proving rigorously (for almost all the coefficients where this is true in the homogeneous case) mass conservation (6) and thus the absence of gelation, a well-known phenomenon in coagulation-fragmentation models [11, 10], where the formal conservation of mass is violated as clusters of infinite size are formed.

In this paper we will work with the global weak solutions constructed in [15] under the assumption

$$\lim_{j \rightarrow +\infty} \frac{a_{i,j}}{j} = \lim_{j \rightarrow +\infty} \frac{B_{i+j} \beta_{i+j,i}}{i+j} = 0, \quad (\text{for fixed } i \geq 1), \quad (7)$$

which were later extended in [20] to the case of  $\Omega = \mathbb{R}^N$ . The notion of solution is the following, which we take from [15]:

**Definition 1.1.** A global weak solution  $c = \{c_i\}_{i \geq 1}$  to (1) – (2) is a sequence of functions  $c_i : [0, +\infty) \times \Omega \rightarrow [0, +\infty)$  such that for each  $T > 0$ ,

$$c_i \in \mathcal{C}([0, T]; L^1(\Omega)), \quad i \geq 1, \quad (8)$$

$$\sum_{j=1}^{\infty} a_{i,j} c_i c_j \in L^1([0, T] \times \Omega), \quad (9)$$

$$\sup_{t \geq 0} \int_{\Omega} \left[ \sum_{i=1}^{\infty} i c_i(t, x) \right] dx \leq \int_{\Omega} \left[ \sum_{i=1}^{\infty} i c_i^0(x) \right] dx, \quad (10)$$

and for each  $i \geq 1$ ,  $c_i$  is a mild solution to the  $i$ -th equation in (1a), that is,

$$c_i(t) = e^{d_i A_1 t} c_i^0 + \int_0^t e^{d_i A_1(t-s)} Q_i[c(s)] ds, \quad t \geq 0, \quad (11)$$

where  $Q_i[c]$  is defined by (2),  $A_1$  denotes the closure in  $L^1(\Omega)$  of the unbounded linear operator  $A$  of  $L^2(\Omega)$  defined by

$$D(A) := \{w \in H^2(\Omega) \mid \nabla w \cdot n = 0 \text{ on } \partial\Omega\}, \quad Aw = \Delta w, \quad (12)$$

and  $e^{d_i A_1 t}$  is the  $C_0$ -semigroup generated by  $d_i A_1$  in  $L^1(\Omega)$ .

The existence result of [15] reads:

**Theorem 1.2** (Laurençot-Mischler). *Assume hypotheses (3) and (7) on the coagulation and fragmentation coefficients. Assume also that*

$$d_i > 0 \quad \text{for all } i \geq 1,$$

*and that the non-negative initial datum has finite mass:*

$$c_i^0 \geq 0 \text{ on } \Omega \quad \text{and} \quad \int_{\Omega} \sum_{i=1}^{\infty} i c_i^0 < +\infty.$$

*Then, there exists a global weak solution to the initial-boundary problem (1) – (2) in the sense of Definition 1.1.*

Under the extra assumptions on the diffusion constants and the initial data

$$0 < \inf_i \{d_i\} =: d, \quad D := \sup_i \{d_i\} < +\infty, \quad (13)$$

$$\sum_{i=1}^{\infty} i c_i^0 \in L^2(\Omega), \quad (14)$$

we are in fact able to prove the following  $L^2$  bound on the mass density  $\rho(t, x) := \sum_{i=1}^{\infty} i c_i(t, x)$ : Denoting by  $\Omega_T$  the cylinder  $[0, T] \times \Omega$ , we have the

**Proposition 1.3.** *Assume that (3), (7), (13) and (14) hold. Then, for all  $T > 0$  the mass  $\rho$  of a weak solution to system (1) – (2) (given by Theorem 1.2) lies in  $L^2(\Omega_T)$  and the following estimate holds:*

$$\|\rho\|_{L^2(\Omega_T)} \leq \left(1 + \frac{\sup_i \{d_i\}}{\inf_i \{d_i\}}\right) T \|\rho(0, \cdot)\|_{L^2(\Omega)}. \quad (15)$$

*Remark 1.4.* Note that the assumption (7) is only included in Proposition 1.3 in order to ensure the existence of a weak solution via Theorem 1.2. Without assumption (7), the bound (15) would still hold for smooth solutions of a truncated version of system (1) – (2) uniformly with respect to the truncation. See [15] for the details of such a truncation.

In addition to Proposition 1.3, we give a new proof of an  $L^1$  bound of the various coagulation and fragmentation terms:

**Proposition 1.5.** *We still assume that (3), (7), (13) and (14) hold. Then, for all  $T > 0$  and  $i \in \mathbb{N}^*$  all the terms  $Q_i^+$ ,  $Q_i^-$ ,  $F_i^+$  and  $F_i^-$  associated to a weak solution to system (1)–(2) (given by Theorem 1.2) lie in  $L^1(\Omega_T)$  with a bound which depends in an explicit way on the coagulation and fragmentation coefficients, the diffusion coefficients, and the initial data  $c_i^0$ .*

*Remark 1.6.* The fact that the terms  $Q_i^+$ ,  $Q_i^-$ ,  $F_i^+$  and  $F_i^-$  associated to a weak solution are in  $L^1(\Omega_T)$  is included in the definition of weak solution; the main content of Proposition 1.5 is the explicit dependence of the bounds on the coefficients and initial data, which can be used to obtain uniform estimates for approximated solutions as we show for instance in section 3. For details on the explicit  $L^1$  bounds we refer to the proof of Proposition 1.5 in section 2.

*Remark 1.7.* The  $L^1$  bounds on  $Q_i^+$ ,  $Q_i^-$ ,  $F_i^+$  and  $F_i^-$  require the assumption (7) only to ensure existence. They would hold at the formal level (that is, for smooth solutions of a truncated system) under the less stringent assumption

$$K_i := \sup_{j \in \mathbb{N}} \frac{B_{i+j} \beta_{i+j,i}}{i+j} < +\infty \quad (i \in \mathbb{N}^*). \quad (16)$$

Note that the above  $L^1$  bound also holds when assumptions (3), (7) are replaced by the assumptions of Theorem 1.2 in [15], but the proof is then much more difficult as it requires an induction on  $i$  which can be removed under our extra assumptions.

In section 3, as a first application of the bounds obtained in Propositions 1.3 and 1.5, we give a very simple proof of existence of weak solutions to

(1)–(2) in dimension  $N = 1$  (that is, the result of Theorem 1.2 in dimension 1) under the additional assumptions (3) and (7).

Our main application of the Propositions 1.3 and 1.5 is however related to the problem of conservation of mass (6), which holds rigorously for solutions to a truncated system (see e.g [15]). Nevertheless, it is an important issue in coagulation-fragmentation theory whether (6) holds for weak solutions of system (1) – (2) itself, or if (6) is replaced by an inequality stating that mass is non-increasing in time. If at some time  $t$ , the identity (6) does not hold any more, we say that gelation occurs, which means from a physical point of view that a macroscopic object has been created.

Our main result in section 4 basically shows that (under the assumptions (3) and (7)) gelation does not occur when the coagulation coefficients  $a_{i,j}$  are at most linear and, moreover, slightly sublinear far off the diagonal  $i = j$ . More precisely, we prove mass conservation under the following condition on the coefficients  $a_{i,j}$ :

**Hypothesis 1.8.** *There is some bounded function  $\theta : [0, +\infty) \rightarrow (0, +\infty)$  such that  $\theta(x) \rightarrow 0$  when  $x \rightarrow +\infty$  and*

$$a_{i,j} \leq (i + j) \theta(j/i) \quad \text{for all } j \geq i. \quad (17)$$

(Or equivalently, by symmetry,

$$a_{i,j} \leq (i + j) \theta(\max\{j/i, i/j\}) \quad \text{for all } i, j \geq 1.)$$

**Theorem 1.9.** *Assume that (3), (7), (13), and (14) hold. Also, assume Hypothesis 1.8. Then, the weak solution to the system (1) given by Theorem 1.2 has a superlinear moment which is bounded on bounded time intervals; this is, there is some increasing function  $C = C(T) > 0$ , and some increasing sequence of positive numbers  $\{\psi_i\}_{i \geq 1}$  with*

$$\lim_{i \rightarrow \infty} \psi_i \rightarrow +\infty \quad (18)$$

such that for all  $T > 0$ ,

$$\int_{\Omega} \sum_{i=1}^{\infty} i \psi_i c_i \leq C(T) \quad \text{for all } t \in [0, T]. \quad (19)$$

As a consequence, under these conditions all weak solutions given by Theorem 1.2 of (1) conserve mass:

$$\int_{\Omega} \rho_0(x) dx = \int_{\Omega} \rho(t, x) dx \quad \text{for all } t \geq 0. \quad (20)$$

*Remark 1.10* (Admissible coagulation coefficients). Let us comment on Hypothesis 1.8. First note that (1.8) includes coefficients of the form

$$a_{i,j} \leq \text{Cst} (i^\alpha j^\beta + i^\beta j^\alpha)$$

for any  $\alpha, \beta > 0$  such that  $\alpha + \beta \leq 1$  (take  $\theta(x) = x^{-\varepsilon}$  for  $\varepsilon > 0$  small enough). It is also satisfied when

$$a_{i,j} \leq \text{Cst} \left( \frac{i}{\phi(i)} + \frac{j}{\phi(j)} \right),$$

where  $x \mapsto \phi(x)$  is any positive strictly increasing function (for  $x$  big enough), which goes to infinity at infinity, and such that  $x \mapsto \frac{x}{\phi(x)}$  is also increasing (take  $\theta(\lambda) = \phi(\lambda)^{-1/2}$ ). All the examples  $\phi = \log(1 + \cdot)$ ,  $\phi = \log(1 + \cdot) \circ \log(1 + \cdot)$ ,  $\dots$ ,  $\phi = \log(1 + \cdot) \circ \dots \circ \log(1 + \cdot)$ ) satisfy this condition. Likewise, condition (17) also holds when (for  $i, j \geq 2$ )

$$a_{ij} \leq \text{Cst} \left( i \frac{R(\log j)}{\log i} + j \frac{R(\log i)}{\log j} \right) \quad (21)$$

for some nondecreasing function  $R$  such that  $x \mapsto R(x)/x$  is nonincreasing and tends to 0 when  $x \rightarrow +\infty$ . Note indeed that when (21) holds,

$$\frac{a_{ij}}{i+j} \leq \frac{1}{1+j/i} \frac{R[\log(j/i) + \log i]}{\log i} + \frac{j/i}{1+j/i} \frac{R[\log i]}{\log(j/i) + \log i}. \quad (22)$$

Then, condition (17) is obtained by distinguishing the cases  $i \geq j/i$  and  $i \leq j/i$  in both terms of the right hand side of (22).

Assumption (21) can even be replaced by

$$a_{ij} \leq \text{Cst} \left( i \frac{R(\log(\log j))}{\log(\log i)} + j \frac{R(\log(\log i))}{\log(\log j)} \right),$$

with the same requirements on  $R$  as previously.

Note however that the linear coefficient  $a_{ij} = i + j$  (or the coefficient  $a_{ij} = \frac{i}{\log i} \log j + \frac{j}{\log j} \log i$ ) does not satisfy hypothesis (1.8), though one would expect that Thm. 1.9 still holds for such coefficients.

Before introducing a generalised coagulation-fragmentation model and thus, a third application of the Propositions 1.3 and 1.5, let us briefly review previous results on existence theory and mass conservation for the coagulation-fragmentation system (1). With some further restrictions on the coefficients as compared to [15], existence of solutions by means of  $L^\infty$

bounds on the  $c_i$  has been proven in [3, 7, 13, 18, 19]. A different technique was used in [1] to prove that equation (1) is well posed, locally in time, and globally in time when the space dimension  $N$  is one, always assuming that the coagulation and fragmentation coefficients are bounded.

In a recent work [14], Hammond and Rezakhanlou considered equation (1) without fragmentation, and gave  $L^\infty$  bounds on moments of the solution (and as a consequence,  $L^\infty$  bounds on the  $c_i$ ). This implies uniqueness and mass conservation for some coagulation coefficients that grow at most linearly as well as an alternative proof of the existence of  $L^\infty$  solutions by a-priori bounds on the  $c_i$ ; for instance, if  $\Omega = \mathbb{R}^N$  and diffusion coefficients  $d_i$  are nonincreasing and satisfying (13) and if moreover

$$\sum_{i=1}^{\infty} i c_i^0 \in L^\infty(\mathbb{R}^N), \quad \sum_{i=1}^{\infty} i^2 c_i^0 \in L^1(\mathbb{R}^N), \quad a_{i,j} \leq C(i+j)$$

for some  $C > 0$  and all  $i, j \geq 1$ , then they show that mass is conserved for all weak solutions of eq. (1) without fragmentation. See [14, Theorems 1.3 and 1.4] and [14, Corollary 1.1] for more details.

In the spatially homogeneous case, mass conservation is known for general data with finite mass and coagulation coefficients including the critical linear case  $a_{i,j} \leq Cst(i+j)$  (see, for instance, [2, 5]).

We finally give a third application of the Propositions 1.3 and 1.5. As mentioned already in the Remarks 1.4 and 1.7, Propositions 1.3 and 1.5 (despite true without restrictions on the coagulation coefficients  $a_{i,j}$  for smooth approximating solutions) do not really improve the theory of existence of weak solutions for the usual models of coagulation-fragmentation like (1) as the full assumption (7) are needed in passing to the limit in the approximating solutions. At best they help provide simpler proofs in particular cases, as done in section 3.

On the other hand, Propositions 1.3 and 1.5 are well suited for the existence theory of more exotic models, for instance, when fragmentation occurs due to binary collisions between clusters. Then, the break-up terms are quadratic, being proportional to the concentration of the two clusters which collide. This leads to coagulation-fragmentation models where all terms in the right hand side are quadratic.

More precisely, we consider that clusters of size  $k$  and  $l$  collide with a rate  $b_{k,l} \geq 0$ , leading to fragmentation. As a consequence, clusters of size  $i < \max\{k, l\}$  are produced, in average, at a rate  $\beta_{i,k,l} \geq 0$  in such a way that the mass is conserved (that is,  $\sum_{i < \max\{k,l\}} i \beta_{i,k,l} = k + l$ ). This leads to the

following system (for  $t \in \mathbb{R}_+$ ,  $x \in \Omega$  a bounded regular open subset of  $\mathbb{R}^N$ ):

$$\begin{aligned} \partial_t c_i - d_i \Delta_x c_i &= \frac{1}{2} \sum_{k+l=i} a_{k,l} c_k c_l - \sum_{k=1}^{\infty} a_{i,k} c_i c_k \\ &+ \frac{1}{2} \sum_{k,l=1}^{\infty} \sum_{i < \max\{k,l\}} b_{k,l} c_k c_l \beta_{i,k,l} - \sum_{k=1}^{\infty} b_{i,k} c_i c_k \quad (i \in \mathbb{N}^*), \end{aligned} \quad (23)$$

together with the initial and boundary conditions (1b), (1c). For this model, the set of assumptions (3) is replaced by

$$a_{i,j} = a_{j,i} \geq 0, \quad (i, j \in \mathbb{N}^*), \quad (24a)$$

$$\beta_{i,k,l} = \beta_{i,l,k} \geq 0, \quad (i, k, l \in \mathbb{N}^*, i < \max\{k, l\}), \quad (24b)$$

$$b_{i,k} = b_{k,i} \geq 0, \quad b_{1,1} = 0, \quad (i, k \in \mathbb{N}^*, i < k), \quad (24c)$$

$$\sum_{i < \max\{k,l\}} i \beta_{i,k,l} = k + l, \quad (k, l \in \mathbb{N}^*). \quad (24d)$$

Because of the quadratic character of the fragmentation terms, the inductive method for the proof of existence devised by Laurençot-Mischler [15] seems difficult to adapt in this case. The method presented in our first application can however be adapted, provided that the dimension is  $N = 1$  and that the following assumptions are made on the coefficients:

**Hypothesis 1.11.** *Assume (24), and suppose that the diffusion coefficients are uniformly bounded above and below (eq. (13)) and that the initial mass lies in  $L^2(\Omega)$  (eq. (14)). In place of (7) we assume further that*

$$\lim_{l \rightarrow \infty} \frac{a_{k,l}}{l} = 0, \quad \lim_{l \rightarrow \infty} \frac{b_{k,l}}{l} = 0, \quad (\text{for fixed } k \in \mathbb{N}^*), \quad (25)$$

$$\lim_{l \rightarrow \infty} \sup_k \left\{ \frac{b_{k,l}}{kl} \beta_{i,k,l} \right\} = 0. \quad (\text{for fixed } i \in \mathbb{N}^*), \quad (26)$$

We define a solution to (23) along the same lines as in Definition 1.1:

**Definition 1.12.** A global weak solution  $c = \{c_i\}_{i \geq 1}$  to (23), the boundary condition (1b) and the initial data (1c) is a sequence of functions  $c_i : [0, +\infty) \times \Omega \rightarrow [0, +\infty)$  such that for each  $T > 0$ ,

$$c_i \in \mathcal{C}([0, T]; L^1(\Omega)), \quad i \geq 1, \quad (27)$$

the four terms on the r.h.s. of (23) are in  $L^1([0, T] \times \Omega)$ ,

$$\sup_{t \geq 0} \int_{\Omega} \left[ \sum_{i=1}^{\infty} i c_i(t, x) \right] dx \leq \int_{\Omega} \left[ \sum_{i=1}^{\infty} i c_i^0(x) \right] dx, \quad (28)$$

and for each  $i \geq 1$ ,  $c_i$  is a mild solution to the  $i$ -th equation in (23), that is,

$$c_i(t) = e^{d_i A_1 t} c_i^0 + \int_0^t e^{d_i A_1 (t-s)} Z_i[c(s)] ds, \quad t \geq 0,$$

where  $Z_i[c]$  represents the right hand side of (23) and  $A_1$ ,  $e^{d_i A_1 t}$  are the same as in Definition 1.1.

We are now able to prove the following theorem:

**Theorem 1.13.** *Under Hypothesis 1.11 on the coefficients and initial data of the equation, and in dimension  $N = 1$ , there exists a global weak solution to eq. (23) satisfying*

$$c_i \in C([0, T], L^1(\Omega)) \cap L^{3-\varepsilon}(\Omega_T) \quad (\text{for all } i \in \mathbb{N}^*, T > 0, \varepsilon > 0),$$

for which the four terms appearing in the right hand side of (23) lie in  $L^1(\Omega_T)$ .

*Remark 1.14.* The method of proof unfortunately does not seem to provide existence in dimensions  $N \geq 2$ . Dimension  $N = 2$  looks in fact critical as it doesn't allow a-priori a bootstrap in the heat equation with right hand side in  $L^1$ . A possible line of proof could follow [12] in the context of reaction-diffusion equations. In higher dimensions  $N \geq 3$ , assuming additionally a detailed balance relation between coagulation and fragmentation, an entropy based duality method as in [9] could be used to define global weak  $L^2$  solutions (see also [16]).

Our paper is built in the following way: Section 2 is devoted to the proof of Propositions 1.3 and 1.5. Then Sections 3, 4, and 5 are each devoted to one of the three applications. In particular, Theorem 1.9 is proven in Section 4 first in a particular case (with a very short proof), and then in complete generality. Theorem 1.13 is proven in Section 5. Finally, an Appendix is devoted to the proof of a Lemma of duality due to M. Pierre and D. Schmitt (cf. [17]), which is the key to Proposition 1.3.

## 2 A new a priori estimate

The solutions given in [15] are constructed by approximating the system (1)–(2) by a truncated system (the procedure consists in setting the coagulation and fragmentation coefficients to zero beyond a given finite size, and smoothing the initial data) for which very regular solutions exist. Then, uniform estimates for the solutions of this approximate system are proven. Finally, it

is shown that these solutions have a subsequence which converges to a solution to the original system. In the proofs below it must be understood that the bounds are obtained for the truncated system (in a uniform way) and then transferred to the weak solution by a passage to the limit: the fact that this transfer can be done (in the case of the total mass) without replacing the equality by an inequality is the heart of our second application.

We begin with the

*Proof of Proposition 1.3.* Using the fact that

$$\partial_t \rho - \Delta(M\rho) = 0, \quad \inf_{i \in \mathbb{N}^*} \{d_i\} \leq M(t, x) := \frac{\sum_{i=1}^{\infty} d_i i c_i}{\sum_{i=1}^{\infty} i c_i} \leq \sup_{i \in \mathbb{N}^*} \{d_i\},$$

we can deduce thanks to a Lemma of duality ([9, Appendix]) that  $\rho \in L^2(\Omega_T)$ , and more precisely that

$$\|\rho\|_{L^2(\Omega_T)} \leq \left(1 + \frac{\sup_i \{d_i\}}{\inf_i \{d_i\}}\right) T \|\rho(0, \cdot)\|_{L^2(\Omega)},$$

for all  $T > 0$ . For the sake of completeness, the Lemma is recalled with its proof in the Appendix (Lemma 6.2).  $\square$

We now turn to the

*Proof of Proposition 1.5.* For  $F_i^-$ , it is clear that

$$F_i^- \leq B_i \rho \in L^2([0, T] \times \Omega) \subseteq L^1([0, T] \times \Omega),$$

thanks to Proposition 1.3. For  $F_i^+$  we use eq. (16) to write

$$F_i^+ \leq \sum_{j=1}^{\infty} \left( \frac{B_{i+j} \beta_{i+j,i}}{i+j} \right) (i+j) c_{i+j} \leq K_i \sum_{j=1}^{\infty} (i+j) c_{i+j} \leq K_i \rho, \quad (29)$$

which is again in  $L^2([0, T] \times \Omega)$ , and hence in  $L^1([0, T] \times \Omega)$ .

For the coagulation terms, we have, since each  $c_i$  is less than  $\rho$ ,

$$Q_i^+ \leq \frac{1}{4} \sum_{j=1}^{i-1} a_{i-j,j} (c_{i-j}^2 + c_j^2) \leq \frac{1}{2} \rho^2 \left( \sum_{j=1}^{i-1} a_{i-j,j} \right), \quad (30)$$

which is in  $L^1([0, T] \times \Omega)$  as  $\rho^2$  is, and the sum only has a finite number of terms. Finally, for  $Q_i^-$  we use the fact that  $Q_i^+$  and  $F_i^+$  are already known

to be integrable: Thus, from eq. (1) integrated over  $[0, T] \times \Omega$ ,

$$\begin{aligned} \int_{\Omega} c_i(T, x) dx + \int_0^T \int_{\Omega} Q_i^-(t, x) dx dt \\ \leq \int_{\Omega} c_i^0(x) dx + \int_0^T \int_{\Omega} Q_i^+(t, x) dx dt + \int_0^T \int_{\Omega} F_i^+(t, x) dx dt. \end{aligned}$$

This proves our result.  $\square$

### 3 First application: a simplified proof of existence of solutions in dimension 1

We begin this section with the following corollary of Proposition 1.5, in the particular case of dimension  $N = 1$ .

**Lemma 3.1.** *Assume that the dimension  $N = 1$ , and that (3), (13), (14) and (16) hold. Then, for all  $T \geq 0$ ,  $i \in \mathbb{N}^*$  the concentrations  $c_i \in L^\infty([0, T] \times \Omega)$  (where  $c_i$  are smooth solutions of a truncated version of (1) – (2), the  $L^\infty$  norm being independent of the truncation).*

*Proof of Lemma 3.1.* We carry out a bootstrap regularity argument. Thanks to Proposition 1.5, we know that (for all  $i \in \mathbb{N}^*$ )

$$(\partial_t - d_i \Delta) c_i \in L^1([0, T] \times \Omega).$$

Using for example the results in [8], this implies that for any  $\delta > 0$ ,

$$c_i \in L^{3-\delta}([0, T] \times \Omega) \quad (i \in \mathbb{N}^*). \quad (31)$$

Now, eq. (31) shows that  $Q_i^+$  is actually more regular: from (the first inequality in) (30),

$$Q_i^+ \in L^{\frac{3}{2}-\frac{\delta}{2}}([0, T] \times \Omega) \quad \text{for all } \delta > 0, i \in \mathbb{N}^*. \quad (32)$$

In addition, we already knew from eq. (29) that (for all  $i \in \mathbb{N}^*$ )

$$F_i^+ \in L^2([0, T] \times \Omega), \quad (33)$$

[for which we do not need to assume that the space dimension  $N$  is 1]. Consequently, omitting the negative terms (for all  $i \in \mathbb{N}^*$ ,  $\delta > 0$ ), we can find  $h_i$  such that

$$(\partial_t - d_i \Delta) c_i \leq h_i \in L^{\frac{3}{2}-\frac{\delta}{2}}([0, T] \times \Omega).$$

As the  $c_i$  are positive, this implies that

$$c_i \in L^p([0, T] \times \Omega) \quad \text{for all } p \in [1, +\infty[, i \in \mathbb{N}^*.$$

Again from (30),

$$Q_i^+ \in L^p([0, T] \times \Omega) \quad \text{for all } p \in [1, +\infty[, i \in \mathbb{N}^*.$$

From this and (33), we can find  $h_i$  such that

$$(\partial_t - d_i \Delta) c_i \leq h_i \in L^2([0, T] \times \Omega),$$

which implies in turn that  $c_i \in L^\infty([0, T] \times \Omega)$  (for all  $i \in \mathbb{N}^*$ ).  $\square$

We now have the possibility to give a short proof of Theorem 1.2 in dimension 1 (and under the extra assumptions (13), (14)). Recall that a proof for any dimension can be found in [15].

*Short proof of Theorem 1.2 in 1D under the assumptions (13) and (14).*

Consider a sequence  $c_i^M$  of (regular) solutions to a truncated version of system (1) – (2). Thanks to Proposition 3.1, we know that for each  $i \in \mathbb{N}^*$ ,  $\sup_M \|c_i^M\|_{L^\infty(\Omega_T)} < +\infty$ . Then (for each  $i \in \mathbb{N}^*$ ) there is a subsequence of the  $(c_i^M)_{M \in \mathbb{N}}$  (which we still denote by  $(c_i^M)_{M \in \mathbb{N}}$ ), and a function  $c_i \in L^\infty(\Omega_T)$ , such that

$$c_i^M \xrightarrow{*} c_i \quad \text{weak-* in } L^\infty(\Omega_T). \quad (34)$$

Using Proposition 1.5, we also see that (for any fixed  $i \in \mathbb{N}^*$ ), the  $L^1(\Omega_T)$  norms of  $C_i^{+,M}$ ,  $C_i^{-,M}$ ,  $F_i^{+,M}$ ,  $F_i^{-,M}$  (the coagulation and fragmentation terms associated to  $\{c_i^M\}$ ) are bounded independently of  $M$ . Using eq. (1a) and the properties of the heat equation, one sees that for each  $i \in \mathbb{N}^*$ , the sequence  $\{c_i^M\}$  lies in a strongly compact subset of  $L^1(\Omega_T)$ . Hence, by renaming our subsequence again, we may assume that

$$c_i^M \rightarrow c_i \quad \text{in } L^1(\Omega_T) \text{ strong, for all } i \in \mathbb{N}^*. \quad (35)$$

In order to prove that  $\{c_i\}$  is indeed a solution to eq. (1) – (2), let us prove that all terms  $F_i^{+,M}$ ,  $F_i^{-,M}$ ,  $C_i^{+,M}$ ,  $C_i^{-,M}$  converge to the corresponding expressions for  $c_i$ , which we denote by  $F_i^+$ ,  $F_i^-$ ,  $C_i^+$ ,  $C_i^-$ , as usual.

1. Positive fragmentation term: for each fixed  $i$ , the sum

$$F_i^{+,M} = \sum_{j=1}^{\infty} B_{i+j} \beta_{i+j,i} c_{i+j}^M$$

converges to  $F_i^+$  in  $L^1(\Omega_T)$  because the tails of the sum converge to 0 uniformly in  $M$  (this is due to hypothesis (7)):

$$\int_0^T \int_{\Omega} \left| \sum_j B_{i+j} \beta_{i+j,i} (c_{i+j}^M - c_{i+j}) \right| dx dt \leq 2 \left( \sup_{j \geq J_0} \left| \frac{B_{i+j} \beta_{i+j,i}}{i+j} \right| \right) \rho + \sup_{j \leq J_0} \|c_{i+j}^M - c_{i+j}\|_{L^1(\Omega_T)}.$$

2. The negative fragmentation term is just a multiple of  $c_i^M$ , so the convergence in  $L^1(\Omega_T)$  is given by (35).
3. For each fixed  $i$ , the positive coagulation term is a finite sum of terms of the form  $a_{i,j} c_i^M c_j^M$ . Thanks to (34) and (35), this converges to  $a_{i,j} c_i c_j$  in  $L^1(\Omega_T)$ .
4. The negative coagulation term is

$$Q_i^{-,M} = c_i^M \sum_{j=1}^{\infty} a_{i,j} c_j^M.$$

Since  $c_i^M$  converges to  $c_i$  weak-\* in  $L^\infty(\Omega_T)$ , it is enough to prove that  $\sum_{j=1}^{\infty} a_{i,j} c_j^M$  converges to  $\sum_{j=1}^{\infty} a_{i,j} c_j$  strongly in  $L^1(\Omega_T)$ . Observing that

$$\int_0^T \int_{\Omega} \left| \sum_j a_{i,j} (c_j^M - c_j) \right| dx dt \leq 2 \left( \sup_{j \geq J_0} \left| \frac{a_{i,j}}{j} \right| \right) \rho + \sup_{j \leq J_0} \|c_j^M - c_j\|_{L^1(\Omega_T)},$$

we see thanks to (7) and (35) that this convergence indeed holds.

□

## 4 Second application: mass conservation

We begin this section with a very short proof of Theorem 1.9 in a particular case in order to show how estimate (15) works. More precisely, we consider the pure coagulation case with  $a_{i,j} = \sqrt{i j}$  and  $B_i = 0$  (no fragmentation), and with initial data satisfy additionally  $\sum_{i=0}^{\infty} i \log i c_i(0, x) dx < +\infty$  (which is slightly more stringent than only assuming finite initial mass).

Then, using the weak formulation (4) with  $\varphi_i = \log(i)$  (and remembering that  $\log(1+x) \leq \text{Cst} \sqrt{x}$ )

$$\begin{aligned} \frac{d}{dt} \int_{\Omega} \sum_{i=1}^{\infty} i \log i c_i dx &= \int_{\Omega} \sum_{i=1}^{\infty} \sum_{j=1}^{\infty} \sqrt{ij} c_i c_j \left( i \log\left(1 + \frac{j}{i}\right) + j \log\left(1 + \frac{i}{j}\right) \right) dx \\ &\leq 2 \int_{\Omega} \sum_{i=1}^{\infty} \sum_{j=1}^{\infty} ij c_i c_j dx \leq 2 \int_{\Omega} \rho(t, x)^2 dx. \end{aligned} \quad (36)$$

As a consequence, we have for all  $T > 0$

$$\int_{\Omega} \sum_{i=0}^{\infty} i \log i c_i(T, x) dx \leq \int_{\Omega} \sum_{i=0}^{\infty} i \log i c_i(0, x) dx + 2 \int_0^T \int_{\Omega} \rho(t, x)^2 dx dt,$$

which ensures the propagation of the moment  $\int \sum_{i=0}^{\infty} i \log i c_i(\cdot, x) dx$ , and therefore gives a rigorous proof of conservation of the mass for weak solutions of the system: no gelation occurs.

Our general result is obtained through a refinement of this argument under hypothesis (1.8). Before giving the proof of Theorem 1.9 we need two technical lemmas, which will substitute the intermediate step in (36).

**Lemma 4.1.** *Let  $\{\mu_i\}_{i \geq 1}$  and  $\{\nu_i\}_{i \geq 1}$  be sequences of positive numbers such that  $\{\mu_i\}$  is bounded,*

$$\sum_{i=1}^{\infty} \mu_i = +\infty \quad \text{and} \quad \lim_{i \rightarrow +\infty} \nu_i = +\infty.$$

*Then we can find a sequence  $\{\xi_i\}_{i \geq 1}$  of nonnegative numbers such that*

$$\begin{aligned} \sum_{i=1}^{\infty} \xi_i &= +\infty, \\ \xi_i \leq \mu_i \quad \text{and} \quad \psi_i := \sum_{j=1}^i \xi_j &\leq \nu_i \quad \text{for all } i \geq 1. \end{aligned}$$

*Proof.* We may assume that  $\nu_i$  is nondecreasing, for otherwise we can consider  $\tilde{\nu}_i := \inf_{j \geq i} \{\nu_j\}$  instead of  $\nu_i$ . Then, in order to find  $\xi_i$  it is enough to define recursively  $\xi_0 := 0$  and, for  $i \geq 1$ ,

$$\xi_i := \begin{cases} \mu_i & \text{if } \mu_i + \sum_{j=0}^{i-1} \xi_j \leq \nu_i, \\ 0 & \text{otherwise.} \end{cases}$$

By construction,  $\xi_i \leq \mu_i$  for all  $i \geq 1$ , and also  $\sum_{j=1}^i \xi_j \leq \nu_i$  for  $i \geq 1$ , as we are assuming  $\{\nu_i\}$  nondecreasing.

To see that  $\{\xi_i\}$  cannot be summable, suppose otherwise that  $\sum_{i=1}^{\infty} \xi_i = S < +\infty$ . Take a bound  $M > 0$  of  $\{\mu_i\}$ , and choose an integer  $k$  such that  $\nu_i \geq S + M$  for all  $i \geq k$ . Then, by definition,

$$\xi_i = \mu_i \quad \text{for all } i \geq k,$$

which implies that  $\{\xi_i\}$  is not summable, as  $\{\mu_i\}$  is not, and gives a contradiction.  $\square$

**Lemma 4.2.** *Assume (17). There is a nondecreasing sequence of positive numbers  $\{\psi_i\}_{i \geq 1}$  such that  $\psi_i \rightarrow +\infty$  when  $i \rightarrow +\infty$ , and*

$$a_{i,j}(\psi_{i+j} - \psi_i) \leq Cj \quad \text{for all } i, j \geq 1, \quad (37)$$

for some constant  $C > 0$ .

In addition, for a given sequence of positive numbers  $\lambda_i$  with  $\lim_{i \rightarrow +\infty} \lambda_i = +\infty$ , we can choose  $\psi_i$  so that  $\psi_i \leq \lambda_i$  for all  $i$ .

*Proof.* First, we may assume that the function  $\theta$  given in Hypothesis 1.8 is nonincreasing on  $[1, +\infty)$ , as we can always take  $\tilde{\theta}(x) := \sup_{y \geq x} \theta(y)$  instead.

We choose a sequence of nonnegative numbers  $\{\xi_i\}$  by applying Lemma 4.1 with

$$\mu_i := \frac{1}{(1+i) \log(1+i)}, \quad (38)$$

$$\nu_i := \min \left\{ \lambda_i, \frac{1}{\theta(\sqrt{i/2})} \right\}. \quad (39)$$

Note that the conditions in Lemma 4.1 are met: the sequence in the right hand side of (38) is not summable, and the right hand side of (39) goes to  $+\infty$  with  $i$ . If we define  $\psi_i := \sum_{j=1}^i \xi_j$ , then the following is given by Lemma 4.1:

$$\xi_i \leq \frac{1}{(1+i) \log(1+i)}, \quad \psi_i \leq \frac{1}{\theta(\sqrt{i/2})}, \quad \psi_i \leq \lambda_i, \quad i \geq 1,$$

$$\lim_{i \rightarrow +\infty} \psi_i = +\infty.$$

These conditions essentially say that  $\psi_i$  grows slower than  $\log \log(i)$ , slower than  $\theta(\sqrt{i/2})^{-1}$ , and slower than  $\lambda_i$ , yet still diverges as  $i \rightarrow +\infty$ .

We can now prove (37) to hold for these  $\{\psi_i\}$  by distinguishing three cases:

1. For any  $i, j \geq 1$ , as  $\log(1+k) \geq 1/2$  for all  $k \geq 1$ ,

$$\psi_{i+j} - \psi_i = \sum_{k=i+1}^{i+j} \xi_k \leq 2 \sum_{k=i+1}^{i+j} \frac{1}{1+k} \leq 2 \log(i+j+1) - 2 \log(i+1) \leq \frac{2j}{i}.$$

Then, in case  $j \leq i$  we use the fact that  $\theta(x) \leq C_\theta$  for some constant  $C_\theta > 0$  and all  $x > 0$  and have

$$a_{i,j}(\psi_{i+j} - \psi_i) \leq 2C_\theta(i+j)\frac{j}{i} \leq 4C_\theta j, \quad \text{for } j \leq i.$$

2. Secondly, for  $i < j \leq i^2$ ,

$$\begin{aligned} \psi_{i+j} - \psi_i &\leq \sum_{k=i+1}^{2i^2} \xi_k \leq \sum_{k=i+1}^{2i^2} \frac{1}{(k+1)\log(k+1)} \\ &\leq \log \log(2i^2+1) - \log \log(i+1) \leq \log \left( \frac{2 \log(\sqrt{3}i)}{\log(i+1)} \right) \leq C_1, \end{aligned}$$

for some number  $C_1 > 0$ . Thus,

$$a_{i,j}(\psi_{i+j} - \psi_i) \leq C_1 C_\theta (i+j) \leq 2C_1 C_\theta j.$$

3. Finally, for  $j > i^2$ ,

$$\psi_{i+j} - \psi_i \leq \psi_{i+j} = \sum_{k=1}^{i+j} \xi_k \leq \frac{1}{\theta(\sqrt{(i+j)/2})} \leq \frac{1}{\theta(\sqrt{j})},$$

and as  $\theta$  is nonincreasing on  $[1, +\infty)$  (we may assume this; see the beginning of this proof), we have for all  $j > i^2$

$$a_{i,j}(\psi_{i+j} - \psi_i) \leq (i+j)\theta(j/i)\frac{1}{\theta(\sqrt{j})} \leq (i+j)\theta(\sqrt{j})\frac{1}{\theta(\sqrt{j})} = i+j \leq 2j.$$

Together, these three cases show (37) for all  $i, j \geq 1$ .  $\square$

Now we are ready to finish the proof of our result on mass conservation:

*Proof of Theorem 1.9.* As remarked above (cf. beginning of section 2), we will prove the estimate (19) for a regular solution to an approximating system, with a constant  $C(T)$  that does not depend on the regularisation. Then, passing to the limit, the result is true for a weak solution thus constructed.

We consider a solution to an approximating system on  $[0, +\infty)$ , which we still denote by  $\{c_i\}_{i \geq 1}$ . Then, by a version of the de la Vallée-Poussin's

Lemma, (see, for instance, Proposition 9.1.1 in [4] or also proof of Lemma 7 in [6]), there exists a nondecreasing sequence of positive numbers  $\{\lambda_i\}_{i \geq 1}$  (independent of the regularisation of the initial data) which diverges as  $i \rightarrow +\infty$ , and such that

$$\int_{\Omega} \sum_{i=1}^{\infty} i \lambda_i c_i^0 dx < +\infty. \quad (40)$$

If we define  $r_i := \int_{\Omega} i c_i^0$ , note that this is just the claim that one can find  $\lambda_i$  as above with  $\sum_i \lambda_i r_i < +\infty$ .

Taking  $\{\psi_i\}$  as given by Lemma 4.2, such that  $\psi_i \leq \lambda_i$  for all  $i \geq 1$ , we have thus  $\int_{\Omega} \sum_{i=1}^{\infty} i \psi_i c_i^0(x) dx < +\infty$ . Then, as integrating over  $\Omega$  makes the diffusion term vanish due to the no-flux boundary conditions, we estimate

$$\frac{d}{dt} \int_{\Omega} \sum_{i=1}^{\infty} i \psi_i c_i dx \leq \frac{1}{2} \int_{\Omega} \sum_{i,j=1}^{\infty} a_{i,j} c_i c_j ((i+j)\psi_{i+j} - i\psi_i - j\psi_j) dx, \quad (41)$$

where we used that the contribution of the fragmentation term is nonpositive, as can be seen from (4) with  $\varphi_i \equiv i \psi_i$ , and the fact that

$$\sum_{j=1}^{i-1} \beta_{i,j} j \psi_j \leq \psi_i \sum_{j=1}^{i-1} \beta_{i,j} j = i \psi_i,$$

as  $\psi_i$  is nondecreasing and (3c) holds. Continuing from (41), by the symmetry of the  $a_{i,j}$ , and using the inequality (37) from Lemma 4.2, we have

$$\frac{d}{dt} \int_{\Omega} \sum_{i=1}^{\infty} i \psi_i c_i dx \leq \int_{\Omega} \sum_{i,j=1}^{\infty} a_{i,j} c_i c_j i (\psi_{i+j} - \psi_i) dx \leq C \int_{\Omega} \rho^2 dx. \quad (42)$$

Thus, Proposition 1.3 showing  $\rho \in L^2(\Omega_T)$  proves that  $\int_{\Omega} \sum_{i=1}^{\infty} i \psi_i c_i dx$  is bounded on bounded time intervals. Mass conservation is a direct consequence of this.  $\square$

*Remark 4.3* (Absence of gelation via tightness). It is interesting to sketch an alternative proof showing conservation of mass via a tightness argument and without establishing superlinear moments. By introducing the superlinear test sequence  $i\phi_k(i)$  with  $\phi_k(i) = \frac{\log i}{\log k} 1_{i < k} + 1_{i \geq k}$  for all  $k \in \mathbb{N}^*$ , we use the weak formulation (4) to see (as above) that the fragmentation part is nonnegative for superlinear test sequences, and use the symmetry of the  $a_{i,j}$

to reduce summation over the indices  $i \geq j \in \mathbb{N}^*$ , which leads to the estimate

$$\begin{aligned} \frac{d}{dt} \int_{\Omega} \sum_{i=1}^{\infty} c_i i \phi_k(i) dx &\leq \int_{\Omega} \sum_{i \geq j} \sum_{j=1}^{\infty} a_{i,j} [i c_i] [c_j] \left( \frac{\log(1 + \frac{i}{j})}{\log(k)} \mathbb{I}_{i < k} \right. \\ &\quad \left. + \frac{j}{i} \left( \frac{\log(1 + \frac{i}{j})}{\log(k)} \mathbb{I}_{i+j < k} + \frac{\log(\frac{k}{j})}{\log(k)} \mathbb{I}_{j < k \leq i+j} \right) \right) dx. \end{aligned}$$

For the first term, we use  $\log(1 + j/i) \leq j/i$ . Then, for the second and third terms, we distinguish further the areas where  $i/j \leq \log(k)$  and  $i/j > \log(k)$ . When  $i/j \leq \log(k)$ , we estimate  $1 + i/j = 1 + i/j \leq 1 + \log(k)$  and  $k/j \leq 1 + i/j \leq 1 + \log(k)$ , respectively. On the other hand, when  $i/j > \log(k)$ , both the second and the third term are bounded by one. Altogether, we get thanks to assumption (1.8), i.e.  $\frac{a_{i,j}}{i} \leq \text{Cst} \theta(i/j)$  for  $i \leq j$ :

$$\begin{aligned} \frac{d}{dt} \int_{\Omega} \sum_{i=1}^{\infty} c_i i \phi_k(i) dx &\leq \left( \frac{1}{\log(k)} + \frac{\log(1 + \log k)}{\log(k)} \right) \sup_{i \geq j \in \mathbb{N}^*} \left\{ \frac{a_{i,j}}{i} \right\} \int_{\Omega} \rho^2 dx \\ &\quad + \int_{\Omega} \sum_{i \geq j} \sum_{j=1}^{\infty} [i c_i] [j c_j] \frac{a_{i,j}}{i} \mathbb{I}_{i/j > \log(k); j < k} dx \\ &\leq \text{Cst} \left( \frac{\log(1 + \log k)}{\log(k)} + \sup_{i/j \geq \log(k)} \theta \left( \frac{i}{j} \right) \right) \int_{\Omega} \rho^2 dx \end{aligned}$$

and the right hand side tends to zero as  $k \rightarrow \infty$ . Hence, using Proposition 1.3 and integrating over a time interval  $[0, T]$ , we get thanks to a tightness argument that the mass is indeed conserved, and no gelation occurs.

## 5 Third Application: Fragmentation due to collisions in dimension 1

*Proof of Theorem 1.13.* We introduce  $(c_i^M)_M$  a sequence of smooth solutions for a truncated version of eq. (23). We first observe that Proposition 1.3 still holds thanks to the duality estimate, that is  $\rho := \sum_i i c_i \in L^2(\Omega_T)$  for all  $T > 0$ . Estimate (30), in which only the coagulation kernel appears, also holds. Moreover, thanks to (24d),

$$\sum_{k,l} \sum_{\max\{k,l\} > i} b_{k,l} c_k c_l \beta_{ikl} \leq \text{Cst}_i \sum_k \sum_l (k+l) c_k c_l \leq \text{Cst}_i \rho^2 \in L^1(\Omega_T).$$

The loss terms

$$\sum_{k=1}^{\infty} a_{i,k} c_i c_k, \quad \sum_{k=1}^{\infty} b_{i,k} c_i c_k$$

lie then in  $L^1(\Omega_T)$  by integration of the equation on  $[0, T] \times \Omega$ .

Using now eq. (23), we see that (for all  $i \in \mathbb{N}^*$ )  $\partial_t c_i^M - d_i \partial_{xx} c_i^M$  belongs to a bounded subset of  $L^1(\Omega_T)$ . As a consequence,  $c_i^M$  belongs (for all  $i \in \mathbb{N}^*$ ) to a compact subset of  $L^{3-\varepsilon}([0, T] \times \Omega)$  for all  $T > 0$  and  $\varepsilon > 0$ . We denote (for all  $i \in \mathbb{N}^*$ ) by  $c_i$  a limit (in  $L^{3-\varepsilon}([0, T] \times \Omega)$  strong) of a subsequence of  $(c_i^M)_{M \in \mathbb{N}}$  (still denoted by  $(c_i^M)_{M \in \mathbb{N}}$ ).

We now pass to the limit in all terms of the r.h.s. of eq. (23). The first term can easily be dealt with, since it consists of a finite sum. Then, we pass to the limit in the second term:

$$\begin{aligned} & \int_0^T \int_{\Omega} \left| \sum_{k=1}^{\infty} a_{i,k} c_i^n c_k^n - \sum_{k=1}^{\infty} a_{i,k} c_i c_k \right| dx dt \\ & \leq \int_0^T \int_{\Omega} \left| \sum_{k=1}^K a_{i,k} c_i^n c_k^n - \sum_{k=1}^K a_{i,k} c_i c_k \right| dx dt + 2 \|\rho\|_{L^2}^2 \sup_{k>K} \left\{ \frac{a_{i,k}}{k} \right\}. \end{aligned}$$

The second part of this expression is small when  $K$  is large enough thanks to assumption (25), (26), while the first part tends to 0 for all given  $K$ .

The fourth term of the r.h.s. of eq. (23) can be treated exactly in the same way. We now turn to the third term:

$$\begin{aligned} & \int_0^T \int_{\Omega} \left| \sum_{k,l=1}^{\infty} \sum_{i < \max\{k,l\}} b_{k,l} c_k^n c_l^n \beta_{i,k,l} - \sum_{k,l=1}^{\infty} \sum_{i < \max\{k,l\}} b_{k,l} c_k c_l \beta_{i,k,l} \right| dx dt \\ & \leq \int_0^T \int_{\Omega} \left| \sum_{k,l=1}^K \sum_{i < \max\{k,l\}}^{k \leq K, l \leq K} b_{k,l} c_k^n c_l^n \beta_{i,k,l} - \sum_{k,l=1}^K \sum_{i < \max\{k,l\}}^{k \leq K, l \leq K} b_{k,l} c_k c_l \beta_{i,k,l} \right| dx dt \\ & \quad + 4 \|\rho\|_{L^2}^2 \sup_{l \geq K} \sup_{k \in \mathbb{N}} \left\{ \frac{b_{k,l}}{kl} \beta_{i,k,l} \right\}. \end{aligned}$$

Once again, the second term is small when  $K$  is large enough thanks to assumption (25), (26), while the first term tends to 0 for all given  $K$ .  $\square$

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## 6 Appendix: A duality lemma

We recall here results from e.g. [17, 9]. We start with the

**Lemma 6.1.** *Assume that  $z : \Omega_T \rightarrow [0, +\infty)$  satisfies*

$$\begin{aligned} \partial_t z + M \Delta z &= -H && \text{on } \Omega, \\ \nabla z \cdot n &= 0 && \text{on } \partial\Omega, \\ z(T, x) &= 0 && \text{on } \Omega, \end{aligned} \tag{43}$$

where  $H \in L^2(\Omega_T)$ , and  $d_1 \geq M \geq d_0 > 0$ . Then,

$$\|z(0, \cdot)\|_{L^2(\Omega)} \leq \left(1 + \frac{d_1}{d_0}\right) T \|H\|_{L^2(\Omega_T)}. \tag{44}$$

*Proof of Lemma 6.1.* Calculating the time derivative of  $\int_{\Omega} |\nabla z|^2$ , or alternatively multiplying eq. (43) by  $\Delta z$  and integrating on  $\Omega$ , we obtain

$$-\frac{1}{2} \frac{d}{dt} \int_{\Omega} |\nabla z|^2 dx + \int_{\Omega} M (\Delta z)^2 dx = \int_{\Omega} -H \Delta z dx,$$

where the boundary condition on  $z$  was used. Integrating on  $[0, T]$  and taking into account that  $z(T, x) = 0$ ,

$$\begin{aligned} \frac{1}{2} \int_{\Omega} |\nabla z(0, \cdot)|^2 dx + \int_{\Omega_T} M (\Delta z)^2 dx dt &= \int_{\Omega_T} H \Delta z dx dt \\ &\leq \|H\|_{L^2(\Omega_T)} \|\Delta z\|_{L^2(\Omega_T)}. \end{aligned} \tag{45}$$

Using that  $M \geq d_0$  we see that  $\int_{\Omega_T} M (\Delta z)^2 \geq d_0 \|\Delta z\|_{L^2(\Omega_T)}^2$ , so (45) implies

$$d_0 \|\Delta z\|_{L^2(\Omega_T)} \leq \|H\|_{L^2(\Omega_T)}.$$

From this and (43) we have

$$\begin{aligned} \|\partial_t z\|_{L^2(\Omega_T)} &\leq \|M\Delta z\|_{L^2(\Omega_T)} + \|H\|_{L^2(\Omega_T)} \\ &\leq d_1 \|\Delta z\|_{L^2(\Omega_T)} + \|H\|_{L^2(\Omega_T)} \leq \left(1 + \frac{d_1}{d_0}\right) \|H\|_{L^2(\Omega_T)}. \end{aligned}$$

Finally,

$$\|z(0, \cdot)\|_{L^2(\Omega)} \leq \int_0^T \|\partial_s z_s\|_{L^2(\Omega)} ds \leq \left(1 + \frac{d_1}{d_0}\right) T \|H\|_{L^2(\Omega_T)}.$$

□

**Lemma 6.2.** *Assume that  $\rho : \Omega_T \rightarrow [0, +\infty)$  and satisfies*

$$\begin{aligned} \partial_t \rho - \Delta(M\rho) &\leq 0 && \text{on } \Omega, \\ \nabla(\rho M) \cdot n &= 0 && \text{on } \partial\Omega, \end{aligned} \tag{46}$$

where  $M : \Omega_T \rightarrow \mathbb{R}$  is a function which satisfies  $d_1 \geq M \geq d_0 > 0$  for some numbers  $d_1, d_0$ . Then,

$$\|\rho\|_{L^2(\Omega_T)} \leq \left(1 + \frac{d_1}{d_0}\right) T \|\rho(0, \cdot)\|_2.$$

*Proof of Lemma 6.2.* Consider the dual problem (43) – (44) for an arbitrary function  $H \in L^2(\Omega_T)$ , with  $H \geq 0$ . Then,  $z \geq 0$ , and integrating by parts in eq. (43), one finds that

$$\begin{aligned} \int_{\Omega_T} \rho H dxdt &= - \int_{\Omega_T} \rho(\partial_t z + M\Delta z) dxdt \\ &= \int_{\Omega_T} z(\partial_t \rho - \Delta(\rho M)) dxdt + \int_{\Omega} \rho(0, \cdot) z(0, \cdot) dxdt \leq \int_{\Omega} \rho(0, \cdot) z(0, \cdot) dxdt, \end{aligned}$$

where we have used eq. (46), eq. (44) and the boundary conditions on  $\rho M$  and  $z$ . Hence, for any nonnegative function  $H \in L^2(\Omega_T)$ ,

$$\int_{\Omega_T} \rho H dxdt \leq \|\rho(0, \cdot)\|_{L^2(\Omega)} \|z(0, \cdot)\|_{L^2(\Omega)},$$

and thanks to Lemma 6.1,

$$\int_{\Omega_T} \rho H dxdt \leq (1 + d_1/d_0) T \|\rho(0, \cdot)\|_{L^2(\Omega)} \|H\|_{L^2(\Omega_T)}.$$

Remembering that  $\rho \geq 0$ , we obtain by duality:

$$\|\rho\|_{L^2(\Omega_T)} \leq (1 + d_1/d_0) T \|\rho(0, \cdot)\|_{L^2(\Omega)}.$$

This proves the lemma. □

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